

# XI. Spin Angular Momentum or Simply "Spin"

- Focus on Electron's Spin Angular Momentum  
[but "spin" can be assigned to particles other than electron]

## ▪ Why "Spin"?

- Need it to understand periodic table [B, C, N, O, F, Ne]

e.g. p orbitals ( $l=1, m_l=1, 0, -1 \Rightarrow 3$  states)

but periodic feature among elements suggested

p orbitals have  $3 \times \underbrace{2} = 6$  states

from spin

- Early experiments asked for it, in addition to orbital AM

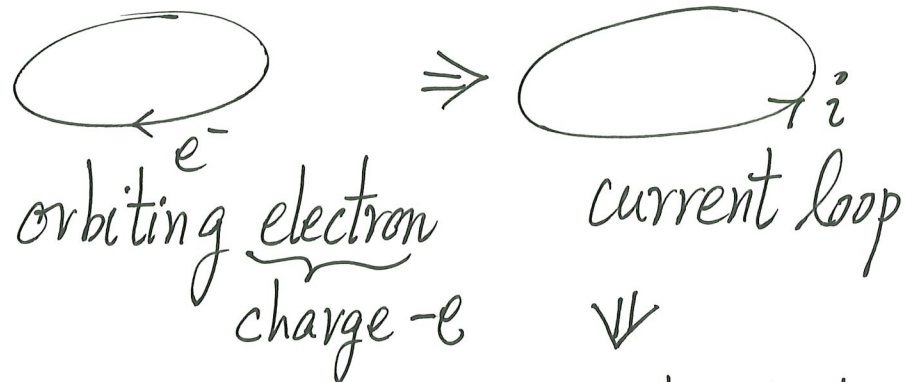
- Need it to explain high-precision hydrogen spectrum

# A. Concept on Experimental Measurement of Angular Momentum

- Discussed Orbital AM  $L^2 \rightarrow l(l+1)\hbar^2$ ,  $l = 0, 1, 2, \dots$   
 $L_z \rightarrow m_l \hbar$ ,  $m_l = \underbrace{l, l-1, \dots, -l}_{(2l+1) \text{ values}}$

- How to do experiments on measuring Orbital AM?

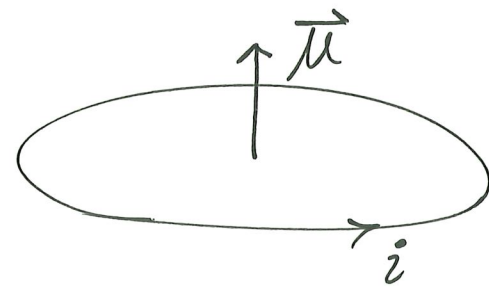
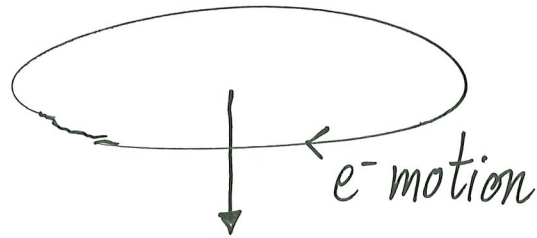
Classical picture



$\vec{A}$  → Magnitude  
 = Area of loop  
 ↘ direction:  $\perp$  plane of loop

$\vec{\mu}$  = magnetic dipole moment  
 $= i\vec{A}$   
 ↙ vector quantity

Thus,



$(\vec{r} \times \vec{p}) \rightarrow \vec{L} = \text{Orbital AM}$

(∵ electron is negatively charged)

∴ Expected  $\vec{\mu}_L \propto -\vec{L}$  (1)

Current  $\rightarrow$

$|i| = e \frac{v}{2\pi r}$

Magnetic dipole moment due to orbital AM (subscript "L")  
[charge passing by a point per unit time]

$|\vec{\mu}_L| = |i| \cdot \pi r^2 = \frac{1}{2} e v r$

Angular Momentum  $\rightarrow$

$|L| = r m_e v$

$\Rightarrow \frac{|\vec{\mu}_L|}{|L|} = \frac{e}{2m_e}$  (2)

Called Gyromagnetic Ratio

Putting (1) & (2) together:

for an electron  $\rightarrow$

$\vec{\mu}_L = \frac{-e}{2m_e} \vec{L}$  (3)

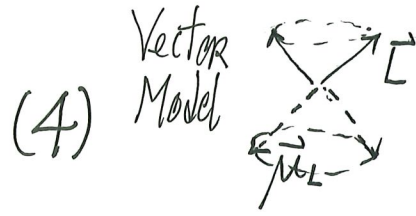
This is "Think Classical"

"Gro Quantum"

$$\hat{\mu}_L = - \frac{e}{2m_e} \hat{L}$$

Operator

a constant



opposite directions

eigenvalues of  $|\mu_L|$

Copy results:

$$|\mu_L| = \frac{e}{2m_e} |\vec{L}| = \frac{e\hbar}{2m_e} \sqrt{l(l+1)}$$

constants  
a combination of  $e, \hbar, m_e$

$$= \mu_B \sqrt{l(l+1)} \quad (5)$$

$\mu_B \equiv \frac{e\hbar}{2m_e} = \text{Bohr Magnetron} = \begin{cases} 9.274 \times 10^{-24} \text{ J/Tesla} \\ 5.79 \times 10^{-5} \text{ eV/Tesla} \end{cases}$

[This is the size of magnetic dipole moment expected of an atom (due to electrons)]

[Note order of magnitude & units]

$(\mu_L)_z = z\text{-component} = \frac{-e}{2m_e} L_z$  takes on  $\frac{-e}{2m_e} m_e \hbar = -\mu_B \cdot m_l$

- But QM deals with Energy (starts with Hamiltonian)
  - Given that  $\vec{L} \rightarrow \vec{\mu}_L$ , how to form an energy term?

$\vec{\mu}_L$  and magnetic field  $\vec{B}$  interact

$$U_{\text{magnetic}} = -\vec{\mu}_L \cdot \vec{B} \quad (6) \quad \text{in an external applied } \vec{B}$$

← an additional term in  $\hat{H}$

- Recall: Atom (spherically symmetric  $U(r)$ )
  - ⇒ No idea about what  $x, y, z$  directions are about!
- Applied  $\vec{B}$  field ⇒  $\vec{B}$  defines a direction
  - call it  $z$ -direction
  - i.e.  $\vec{B} = B \hat{z}$  (no loss of generality)

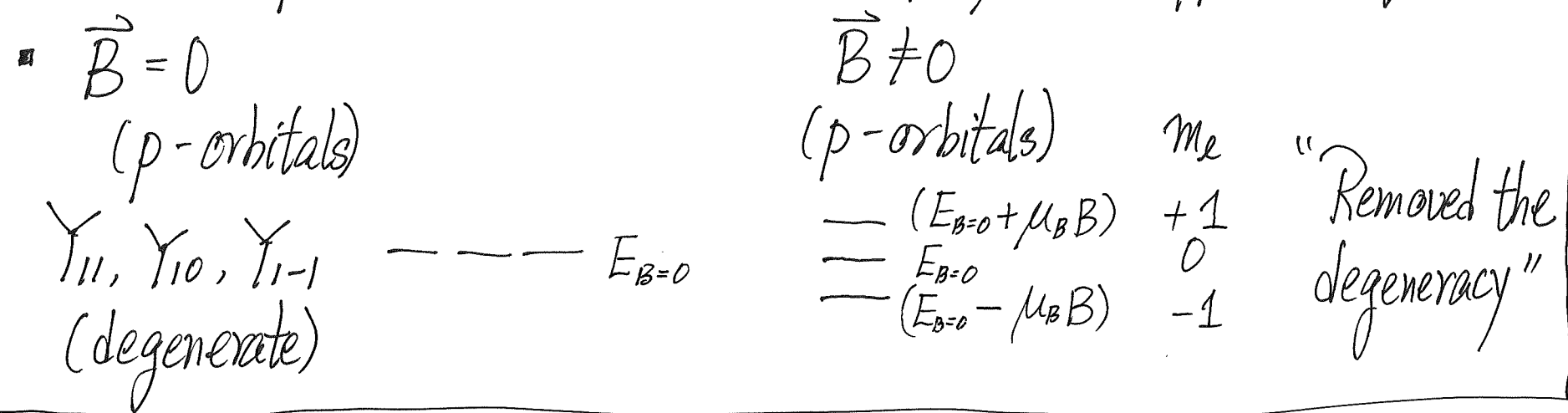
for a state of quantum #  $m_l$  XI-6

$$\therefore U_{\text{magnetic}} = -\vec{\mu}_L \cdot \vec{B} = -(\mu_L)_z B = + \frac{e\hbar}{2m_e} B m_l = \mu_B \cdot B \cdot m_l$$

[ $B \sim$  few Tesla,  $\mu_B \sim 10^{-5}$  eV/Tesla,  $U_{\text{magnetic}} \sim 10^{-4} - 10^{-5}$  eV tiny but detectable]

### Take-Home Message

- electron has charge (-e)  $\rightarrow \vec{L}$  leads to  $\vec{\mu}_L = -\frac{e}{2m_e} \vec{L}$
- $\vec{\mu}_L$  interacts with  $\vec{B} \Rightarrow U_{\text{magnetic}} = -\vec{\mu}_L \cdot \vec{B} = -(\mu_L)_z B$
- To do experiments on orbital AM, play with applied  $\vec{B}$  field



(7)

- This is the physics behind the Zeeman Effect,

one spectral line ( $\vec{B}=0$ ) splits into several ( $\vec{B}\neq 0$ )

(1902 Nobel Prize)

- Note that  $m_l = -1$  has lowest energy

$\vec{\mu}_L$  tends to be "aligned" with  $\vec{B}$  ( $(\mu_L)_z = -\mu_B m_l = +\mu_B$ )

and thus  $\vec{L}$  tends to be "anti-aligned" with  $\vec{B}$  ( $L_z = m_l \hbar = -\hbar$ )

- Important for discussion on Spin

s-orbital (e.g. H-atom in ground state)

$\Rightarrow l=0 \Rightarrow m_l=0$  only

$\Rightarrow$  Don't expect to observe any effect in  $\vec{B}$ -field  
if there is only orbital AM

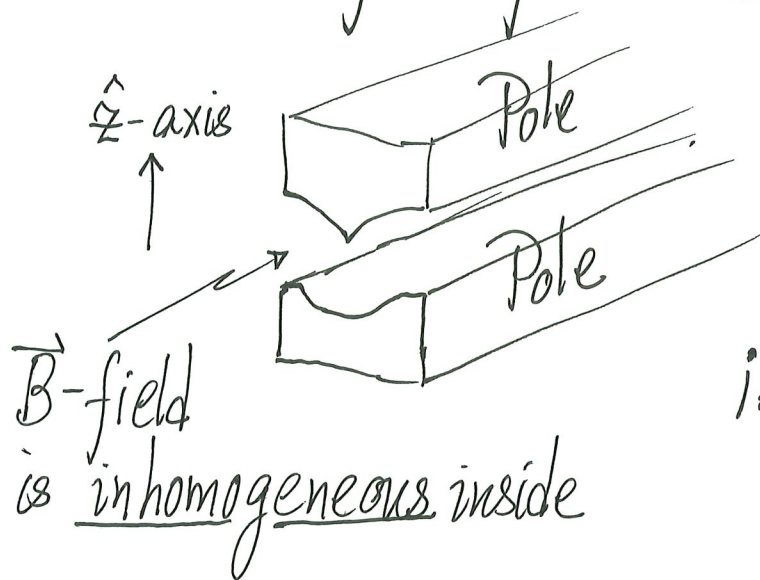
(8)

But this is NOT what experiments showed! (See Sec. C, Stern-Gerlach Expt)

## B. Inhomogeneous Magnetic Field exerts force on magnetic dipole moment

- An example of inhomogeneous (non-uniform) magnetic field

[Uniform field:  $\vec{B}(x,y,z) = \vec{B}$  (same  $\vec{B}$  at different places)]



← Special designed magnet  
for giving inhomogeneous field

i.e.  $\vec{B}(\vec{r})$

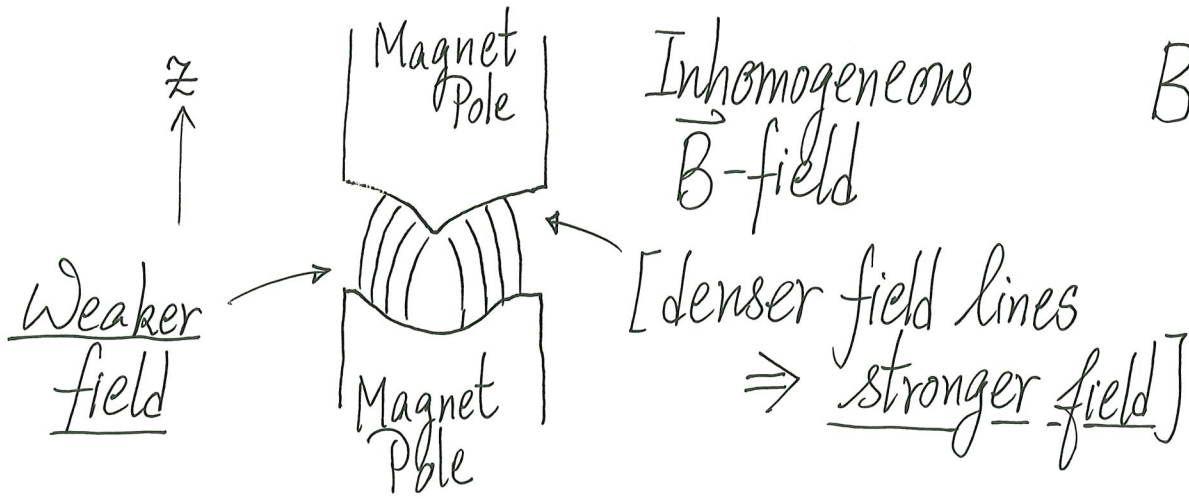
$\vec{B}$ -field  
is inhomogeneous inside

- So what?

- A magnetic dipole moment  $\vec{\mu}$  feels a force in an inhomogeneous  $\vec{B}$ -field

more than a torque





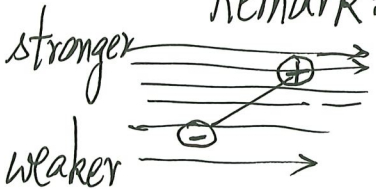
$B(z) \Rightarrow$  Non-uniform  
 ( $B$  is different at different places)

$\therefore$  there is a  $\frac{dB}{dz} \neq 0$

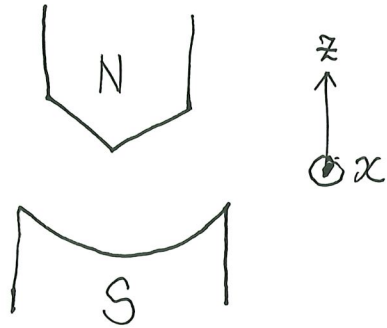
$$\begin{aligned} \vec{F} &= \text{Force}^\dagger \text{ on } \vec{\mu} \text{ due to inhomogeneous } \vec{B}\text{-field} \\ &= -\vec{\nabla}(-\vec{\mu} \cdot \vec{B}) \quad (-\vec{\nabla}(\text{potential energy}) = \text{force}) \\ &= (\vec{\mu} \cdot \vec{\nabla}) \vec{B} \end{aligned}$$

$\therefore$  As long as  $\vec{B}$  varies, there is a force on  $\vec{\mu}$ .

<sup>†</sup>Remark: There is an electric dipole moment in an inhomogeneous  $\vec{E}$ -field analogy.  
 net force on  $\vec{p}_{\text{electric}} \sim (\vec{p}_{\text{el}} \cdot \vec{\nabla}) \vec{E}$



# Key Point



As atom travels down the inhomogeneous field (into the page), it experiences a force

$$F_z = \underbrace{\mu_z}_{\substack{\text{pushing atom} \\ \text{up or down}}} \underbrace{\left( \frac{\partial B}{\partial z} \right)}_{\substack{\text{control by design of magnet} \\ \text{(a constant)}}}$$

Note:  $\frac{\partial B_z}{\partial z}$  more precisely

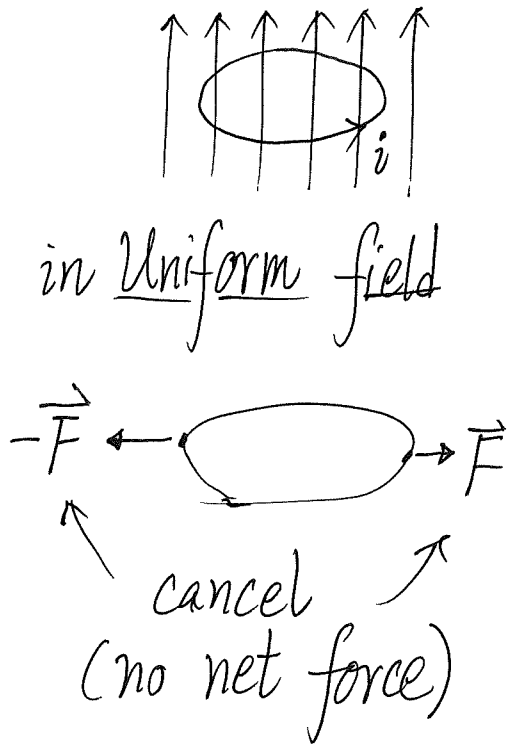
key quantity

z-component of Magnetic dipole moment

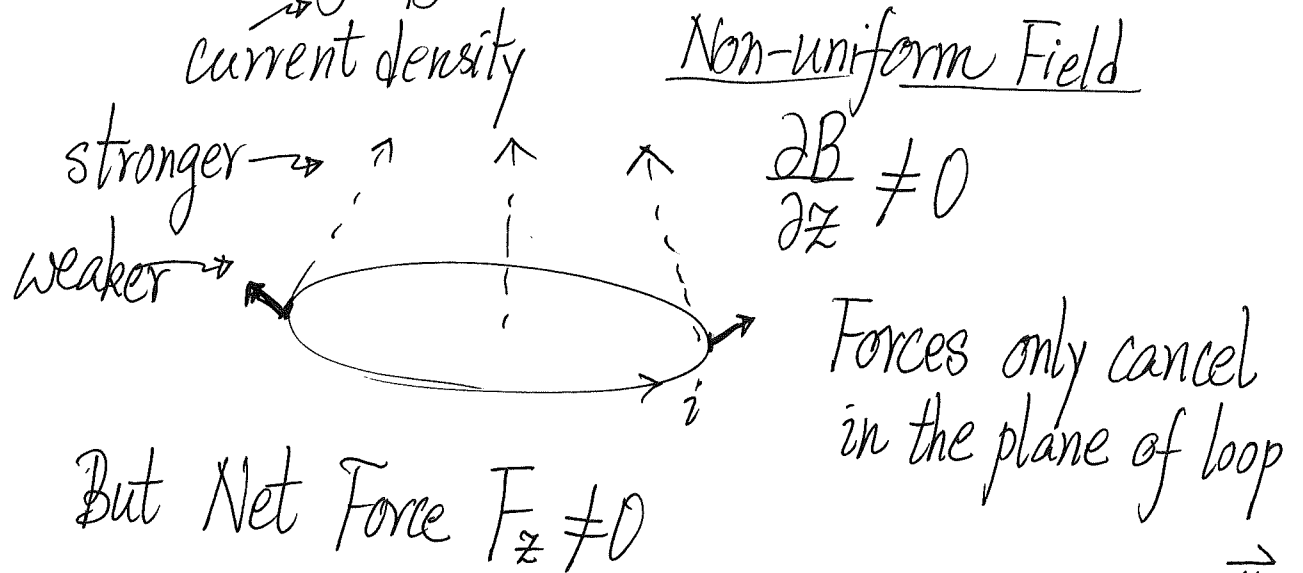
along  $+\hat{z}$  direction  $\Rightarrow F_z$  pushes atom up  
 along  $-\hat{z}$  direction  $\Rightarrow F_z$  pushes atom down

[if atom carries a magnetic dipole moment]

Classical picture:  $\vec{\mu} \approx$  current loop



[Recall:  $(-e)\vec{v} \times \vec{B} =$  force on moving charge]  
 $\sim \vec{J} \times \vec{B}$   
 current density



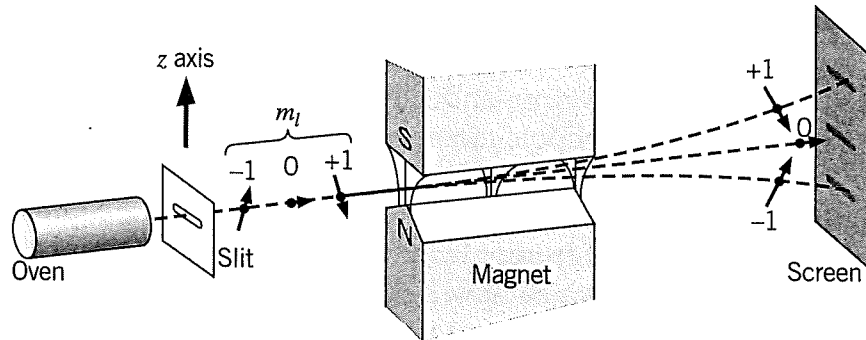
$\Rightarrow$  Net Force in direction of  $\frac{\partial B}{\partial z}$  for  $\vec{\mu}$

For  Net Force in opposite direction

$\therefore F_z \propto \mu_z \frac{\partial B}{\partial z} \Rightarrow$  different  $\mu_z$  experiences different forces

Example: Hydrogen atoms (all excited to 2p states)

- Only knowledge of orbital AM ( $l=1, m_l=1, 0, -1$ )
  - Passing atoms through inhomogeneous  $\vec{B}$ -field  $L_z = \hbar, 0, -\hbar$   
 $(\mu_L)_z = -\mu_B, 0, +\mu_B$  ( $m_l \cdot \mu_B$ )  
 [assuming all stay in 2p]  $\uparrow$  (No force)  $\uparrow$   
 opposite forces
- We would expect to see



Schematic diagram of Stern-Gerlach experiment. A beam of atoms from an oven passes through a slit and then enters a region where there is a nonuniform magnetic field. Atoms with their magnetic dipole moments in opposite directions experience forces in opposite directions.

[Taken from Krane, "Modern Physics"]

Key Point

- Different  $\mu_z$ : different forces
- This experimental set up is called the Stern-Gerlach Experiment (SGE)

Remarks:

- Ideas → electric dipole moment experiences a force in inhomogeneous electric field  
(e.g. atoms seeking strong field locations, polarized spheres seeking strong field locations)
- magnetic dipole moment experiences a force in inhomogeneous magnetic field

are at the heart of many modern physics experiments!  
Optical tweezers, Optical lattice (forming controllable artificial crystals),  
electrorheological / magnetorheological fluids